





# Hadronic interactions in Angantyr

#### Marius Utheim in collaboration with Ilkka Helenius

Wuppertal, January 22nd, 2024



Background

Modelling

Results

Summary and outlook

Background

# Outline

Background

Modelling

Results

Summary and outlook

# Pythia and Angantyr

 $\ensuremath{\operatorname{PYTHIA}}$  is a general-purpose event generator.

- For information about PYTHIA itself, see Torbjörn's talk from earlier
- In this talk, I focus on ANGANTYR, the module for heavy ions.



Current objective: To implement hadron-ion interactions for generic hadrons.

One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.

- One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.
- ► Another use case is for the VMD part of the photon wave function.

- One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.
- Another use case is for the VMD part of the photon wave function.
- ▶ Also relevant for direct comparisons, e.g. to NA61/SHINE  $\pi^-$  C data.

- One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.
- ► Another use case is for the VMD part of the photon wave function.
- ▶ Also relevant for direct comparisons, e.g. to NA61/SHINE  $\pi^-$  C data.
- Technical requirement: change beam types and energies on an event-by-event basis.

- One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.
- ► Another use case is for the VMD part of the photon wave function.
- ▶ Also relevant for direct comparisons, e.g. to NA61/SHINE  $\pi^-$  C data.
- Technical requirement: change beam types and energies on an event-by-event basis.

Current objective: To implement hadron-ion interactions for generic hadrons.

- One main use case is hadronic cascades, e.g. using PYTHIA as the interaction model in CORSIKA 8.
- ► Another use case is for the VMD part of the photon wave function.
- ▶ Also relevant for direct comparisons, e.g. to NA61/SHINE  $\pi^-$  C data.
- Technical requirement: change beam types and energies on an event-by-event basis.

Hadron-proton interactions already exist in PYTHIA [arXiv:2108.03481]. In this talk, I will present this framework, then introduce the changes we have made to extend this to ANGANTYR. Finally we look at some results, comparing to HERA and LHC data.

# Cosmic rays

















▶ The direct part is straightforward to model in ANGANTYR: the photon simply scatters off a single nucleon. At high Q<sup>2</sup>, this corresponds to DIS.



- The direct part is straightforward to model in ANGANTYR: the photon simply scatters off a single nucleon. At high  $Q^2$ , this corresponds to DIS.
- The anomalous part is more complicated. The q and  $\bar{q}$  can interact with different nucleons in A.



- ► The direct part is straightforward to model in ANGANTYR: the photon simply scatters off a single nucleon. At high Q<sup>2</sup>, this corresponds to DIS.
- The anomalous part is more complicated. The q and  $\bar{q}$  can interact with different nucleons in A.
- ▶ The VMD part can be described as a *hA* interaction, analogous to *pA*. This is the component with highest multiplicity, due to MPIs and multiple subcollisions.

# Outline

Background

#### Modelling

Results

Summary and outlook

Several models for total cross sections are available in Pythia. The most generic is the Donnachie-Landshoff model, which is available for most hadron-nucleon combinations:

$$\sigma_{AB}(s) = X^{AB}s^{\epsilon} + Y^{AB}s^{-\eta}$$

Several models for total cross sections are available in Pythia. The most generic is the Donnachie-Landshoff model, which is available for most hadron-nucleon combinations:

$$\sigma_{AB}(s) = X^{AB}s^{\epsilon} + Y^{AB}s^{-\eta}$$

Elastic and diffractive cross sections are based on parameterizations by SaS, e.g.



Several models for total cross sections are available in Pythia. The most generic is the Donnachie-Landshoff model, which is available for most hadron-nucleon combinations:

$$\sigma_{AB}(s) = X^{AB}s^{\epsilon} + Y^{AB}s^{-\eta}$$

Elastic and diffractive cross sections are based on parameterizations by SaS, e.g.



• If  $\beta_{B\mathbb{P}}(t) = \beta_{B\mathbb{P}} \exp(b_B t)$ , then with suitable normalization,  $X^{AB} = \beta_{A\mathbb{P}}(0)\beta_{B\mathbb{P}}(0)$ 

Several models for total cross sections are available in Pythia. The most generic is the Donnachie-Landshoff model, which is available for most hadron-nucleon combinations:

$$\sigma_{AB}(s) = X^{AB}s^{\epsilon} + Y^{AB}s^{-\eta}$$

Elastic and diffractive cross sections are based on parameterizations by SaS, e.g.



• If  $\beta_{B\mathbb{P}}(t) = \beta_{B\mathbb{P}} \exp(b_B t)$ , then with suitable normalization,  $X^{AB} = \beta_{A\mathbb{P}}(0)\beta_{B\mathbb{P}}(0)$ 

►  $B_{XB} = 2b_B + 2\alpha'_{\mathbb{P}} \log(s/M_X^2)$  with b = 1.4 for mesons and 2.3 for baryons

Several models for total cross sections are available in Pythia. The most generic is the Donnachie-Landshoff model, which is available for most hadron-nucleon combinations:

$$\sigma_{AB}(s) = X^{AB}s^{\epsilon} + Y^{AB}s^{-\eta}$$

Elastic and diffractive cross sections are based on parameterizations by SaS, e.g.



• If  $\beta_{B\mathbb{P}}(t) = \beta_{B\mathbb{P}} \exp(b_B t)$ , then with suitable normalization,  $X^{AB} = \beta_{A\mathbb{P}}(0)\beta_{B\mathbb{P}}(0)$ 

►  $B_{XB} = 2b_B + 2\alpha'_{\mathbb{P}} \log(s/M_X^2)$  with b = 1.4 for mesons and 2.3 for baryons

▶ *F*<sub>SD</sub> is a fudge factor (out of scope for this talk)

#### Parton distribution functions

PDFs determine the contents of a hadron. For protons, detailed PDFs based on global fits exist, the Pythia default being NNPDF2.3 QCD+QED LO (with  $\alpha_S = 0.130$ ).

#### Parton distribution functions

PDFs determine the contents of a hadron. For protons, detailed PDFs based on global fits exist, the Pythia default being NNPDF2.3 QCD+QED LO (with  $\alpha_S = 0.130$ ).

For other species, very little data exists, and we base our valence distributions on an ansatz by Glück, Reya et al. [arXiv:hep-ph/9806404]:

$$f(x, Q_0^2 = 0.26 \text{ GeV}^2) = Nx^a(1-x)^b(1+A\sqrt{x}+Bx)$$

and evolve to higher scales using the QCDNUM program. The parameters are fixed by flavour- and momentum sum relations, and some heuristic guesses. In particular, heavier valence quarks should have larger x, as they must all have similar velocities in order for the hadron to stay intact.

## Parton distribution functions



 $\triangleright$   $\langle x \rangle$  is higher for heavy valence content (solid lines), and correspondingly lower for light content (dashed lines).

In a nutshell, ANGANTYR sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using PYTHIA.

In a nutshell,  $\rm Angantyr$  sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using  $\rm PythiA.$ 

 Nuclear geometry is given by Glauber model.
 Each subcollision is assigned a type (absorptive, diffractive, elastic) based on the impact parameter b<sub>NN</sub>.



In a nutshell,  $\rm Angantyr$  sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using  $\rm Pythia.$ 

- Nuclear geometry is given by Glauber model.
   Each subcollision is assigned a type (absorptive, diffractive, elastic) based on the impact parameter b<sub>NN</sub>.
- Geometry includes nuclear fluctuations.



In a nutshell, ANGANTYR sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using PYTHIA.

- Nuclear geometry is given by Glauber model.
   Each subcollision is assigned a type (absorptive, diffractive, elastic) based on the impact parameter b<sub>NN</sub>.
- Geometry includes nuclear fluctuations.
- Perform absorptive subcollisions with smallest b<sub>NN</sub> first. Generate events to parton level.



In a nutshell,  $\rm Angantyr$  sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using  $\rm Pythia.$ 

- Nuclear geometry is given by Glauber model.
   Each subcollision is assigned a type (absorptive, diffractive, elastic) based on the impact parameter b<sub>NN</sub>.
- Geometry includes nuclear fluctuations.
- Perform absorptive subcollisions with smallest b<sub>NN</sub> first. Generate events to parton level.
- Secondary absorptive collisions are modelled like diffractive interactions.



In a nutshell, ANGANTYR sets up the spatial configuration of each nucleus, then proceeds by simulating individual nucleon-nucleon interactions using PYTHIA.

- Nuclear geometry is given by Glauber model.
   Each subcollision is assigned a type (absorptive, diffractive, elastic) based on the impact parameter b<sub>NN</sub>.
- Geometry includes nuclear fluctuations.
- Perform absorptive subcollisions with smallest b<sub>NN</sub> first. Generate events to parton level.
- Secondary absorptive collisions are modelled like diffractive interactions.
- Combine partons from all subevents, then do color reconnection, string interactions, string hadronization, etc.



# Hadronic fluctuations

Hadron-hadron interactions already exist in PYTHIA, including their PDFs and cross sections. The only non-trivial hadronic effect left to consider is hadronic fluctuations.

## Hadronic fluctuations

Hadron-hadron interactions already exist in PYTHIA, including their PDFs and cross sections. The only non-trivial hadronic effect left to consider is hadronic fluctuations. In ANGANTYR, the fluctuations are controlled by three parameters. The size of the fluctuations fix the cross sections. The fluctuation parameters are fitted with a genetic algorithm to reproduce these cross sections.

# Hadronic fluctuations

Hadron-hadron interactions already exist in PYTHIA, including their PDFs and cross sections. The only non-trivial hadronic effect left to consider is hadronic fluctuations. In ANGANTYR, the fluctuations are controlled by three parameters. The size of the fluctuations fix the cross sections. The fluctuation parameters are fitted with a genetic algorithm to reproduce these cross sections.

Applying this procedure to the asymmetric meson-nucleon case gives mediocre results. The fitting targets are insensitive to asymmetric fluctuations. Furthermore, the model gives unphysically large fluctuations for  $J/\psi$ , which is expected to have a small wavefunction. Clearly there is room for improvement.

# Model test: Rapidity spectra at 5.02 TeV



▶ For heavier mesons, the rapidity spectrum is pushed in the meson-going direction.

## Model test: Rapidity spectra at 5.02 TeV



For heavier mesons, the rapidity spectrum is pushed in the meson-going direction.
 Relation between p and ρ<sup>0</sup> is unexpected – it is a model uncertainty due to PDFs.

## Model test: Rapidity spectra at 5.02 TeV



For heavier mesons, the rapidity spectrum is pushed in the meson-going direction.
 Relation between p and ρ<sup>0</sup> is unexpected – it is a model uncertainty due to PDFs.
 pPb has more subcollisions, and is thus pushed harder in the ion-going direction.

## Model test: Rapidity spectra at 5.02 TeV



For heavier mesons, the rapidity spectrum is pushed in the meson-going direction.
 Relation between p and ρ<sup>0</sup> is unexpected – it is a model uncertainty due to PDFs.
 pPb has more subcollisions, and is thus pushed harder in the ion-going direction.
 Due to fluctuations and impact parameter sampling, J/ψ gets some events with extremely high weight.

# Model test: Multiplicities at 5.02 TeV



Bimodal peaks are due to the presence or absence of an absorptive subcollision.

# Model test: Multiplicities at 5.02 TeV



Bimodal peaks are due to the presence or absence of an absorptive subcollision.

Long proton tail is driven by larger cross section and more subcollisions.

# Model test: Multiplicities at 5.02 TeV



Bimodal peaks are due to the presence or absence of an absorptive subcollision.

- ► Long proton tail is driven by larger cross section and more subcollisions.
- Heavier mesons produce fewer subcollisions, but each subcollision produces more particles, leading to a non-trivial progression from  $\rho^0$  to  $\phi$  to  $J/\psi$ .

## Outline

Background

Modelling

Results

Summary and outlook

# $\pi^- C$ at NA61/SHINE [arXiv:2209.10561v1]



▶ The different colours refer to different values of the  $p_{0,\perp}^{\text{ref}}$  parameter, which represents a saturation scale in MPI evolution.

# $\pi^- C$ at NA61/SHINE [arXiv:2209.10561v1]



- ▶ The different colours refer to different values of the  $p_{0,\perp}^{\text{ref}}$  parameter, which represents a saturation scale in MPI evolution.
- ► ANGANTYR shows good agreement in meson spectra.

# $\pi^- C$ at NA61/SHINE [arXiv:2209.10561v1]



- ▶ The different colours refer to different values of the  $p_{0,\perp}^{\text{ref}}$  parameter, which represents a saturation scale in MPI evolution.
- ► ANGANTYR shows good agreement in meson spectra.
- Baryons are less well described.

# $\pi^- C$ at NA61/SHINE [arXiv:2209.10561v1]



- ▶ The different colours refer to different values of the  $p_{0,\perp}^{\text{ref}}$  parameter, which represents a saturation scale in MPI evolution.
- ► ANGANTYR shows good agreement in meson spectra.
- Baryons are less well described.
- Low energy framework is not applied here.

# $\gamma p$ at HERA [arXiv:2106.12377]



▶ The  $p_{0,1}^{\text{ref}}$  variation gives a sense of the model uncertainty.

### $\gamma p$ at HERA [arXiv:2106.12377]



• The  $p_{0,1}^{\text{ref}}$  variation gives a sense of the model uncertainty.

▶ The shift due to changing  $p_{0,\perp}^{\text{ref}}$  is larger on average in the full photoproduction than in just the VMD component.

# ATLAS $\gamma$ + Pb multiplicities [arXiv:2101.10771]



> The ATLAS data is not corrected for the limited efficiency, estimated to  $\sim 80$  %.

# ATLAS $\gamma$ + Pb multiplicities [arXiv:2101.10771]



The ATLAS data is not corrected for the limited efficiency, estimated to ~ 80 %.
 Qualitatively speaking, the shift from γp to γPb is consistent with data.

# ATLAS $\gamma$ + Pb multiplicities [arXiv:2101.10771]



- ▶ The ATLAS data is not corrected for the limited efficiency, estimated to  $\sim 80$  %.
- Qualitatively speaking, the shift from  $\gamma p$  to  $\gamma Pb$  is consistent with data.
- In γp, the VMD component has less average multiplicity than in full photoproduction. This could be the other way around for γPb.

#### ATLAS eta spectrum [arXiv:2101.10771]



Again, we cannot make a direct comparison, but the fit is still good when accounting for the limited efficiency in the multiplicity cut.

# Outline

Background

Modelling

Results

Summary and outlook

In this project, we have implemented hadron-ion collisions for all hadron species.

Primary applications are to cosmic rays and photo-induced processes.

- > Primary applications are to cosmic rays and photo-induced processes.
- ▶ Relies on existing PYTHIA hadron-hadron framework. The main non-trivial new physics feature is hadron size fluctuations. The current model has some flaws, particularly noticeable for  $J/\psi$ .

- Primary applications are to cosmic rays and photo-induced processes.
- ► Relies on existing PYTHIA hadron-hadron framework. The main non-trivial new physics feature is hadron size fluctuations. The current model has some flaws, particularly noticeable for  $J/\psi$ .
- Another significant source of uncertainty is in hadron PDFs.

- Primary applications are to cosmic rays and photo-induced processes.
- ► Relies on existing PYTHIA hadron-hadron framework. The main non-trivial new physics feature is hadron size fluctuations. The current model has some flaws, particularly noticeable for  $J/\psi$ .
- Another significant source of uncertainty is in hadron PDFs.
- Room for other details, such as low-energy modelling.

- > Primary applications are to cosmic rays and photo-induced processes.
- ▶ Relies on existing PYTHIA hadron-hadron framework. The main non-trivial new physics feature is hadron size fluctuations. The current model has some flaws, particularly noticeable for  $J/\psi$ .
- Another significant source of uncertainty is in hadron PDFs.
- Room for other details, such as low-energy modelling.
- The work also includes technical features, in particular energy and beam switching.

- > Primary applications are to cosmic rays and photo-induced processes.
- ► Relies on existing PYTHIA hadron-hadron framework. The main non-trivial new physics feature is hadron size fluctuations. The current model has some flaws, particularly noticeable for  $J/\psi$ .
- Another significant source of uncertainty is in hadron PDFs.
- ▶ Room for other details, such as low-energy modelling.
- The work also includes technical features, in particular energy and beam switching.
- Our model shows a good agreement with data from NA61/SHINE, HERA, and ATLAS UPCs.

Backup slides

# Outline

Backup slides

#### Photon flux [arXiv:1901.05261]



Backup slides

# Photon wavefunction details [arXiv:hep-ph/9403393]







$$|\gamma\rangle = c_{\text{bare}} |\gamma_{\text{bare}}\rangle + \sum_{q} c_{q} |q\bar{q}\rangle + \sum_{V=\rho^{0},\omega,\phi,J/\psi} c_{V} |V\rangle$$

$$c_V = \frac{4\pi\alpha_{EM}}{f_V^2}$$

| V        | $f_V^2/4\pi$ |
|----------|--------------|
| $ ho^0$  | 2.20         |
| $\omega$ | 23.6         |
| $\phi$   | 18.4         |
| $J/\psi$ | 11.5         |

# $p_{0,\perp}^{ m ref}$ variations

